# Eigenvalue problems with unbalanced growth: Nonlinear patterns and standing wave solutions 

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#### Abstract

We consider two classes of nonlinear eigenvalue problems with double-phase energy and lack of compactness. We establish existence and non-existence results and related properties of solutions. Our analysis combines variational methods with the generalized Pucci-Serrin maximum principle.


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## 1. Introduction

In this paper, motivated by recent advances in the study of nonlinear problems with unbalanced growth, we are interested in the mathematical analysis of standing wave solutions of some classes of Dirichlet boundary value problems driven by nonhomogeneous differential operators of the type

$$
\begin{equation*}
\operatorname{div}\left[\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u\right], \tag{1}
\end{equation*}
$$

where $\phi \in C^{1}\left(\mathbb{R}_{+}, \mathbb{R}_{+}\right)$has a different growth near zero and at infinity. Such a behaviour occurs, for instance, if $\phi(t)=2(\sqrt{1+t}-1)$, which corresponds to the prescribed mean curvature differential operator (capillary

[^0]surface operator), namely
$$
\operatorname{div}\left(\frac{\nabla u}{\sqrt{1+|\nabla u|^{2}}}\right)
$$

More generally, $\phi(t)$ can behave like $t^{q / 2}$ for small $t$ and like $t^{p / 2}$ for large $t$, where $1<p<q$. Such a growth is fulfilled if

$$
\phi(t)=\frac{2}{p}\left[\left(1+t^{q / 2}\right)^{p / q}-1\right],
$$

which generates the differential operator

$$
\operatorname{div}\left(\left(1+|\nabla u|^{q}\right)^{(p-q) / q}|\nabla u|^{q-2} \nabla u\right) .
$$

A case intensively studied in recent years corresponds to

$$
\phi(t)= \begin{cases}\frac{2}{q} t^{q / 2} & \text { if } t<1 \\ \frac{2}{p} t^{p / 2}-\frac{2(q-p)}{p q} & \text { if } t \geqslant 1\end{cases}
$$

It follows that

$$
\phi\left(|\nabla u|^{2}\right) \simeq \begin{cases}|\nabla u|^{p}, & \text { if }|\nabla u| \gg 1 ; \\ |\nabla u|^{q}, & \text { if }|\nabla u| \ll 1 .\end{cases}
$$

This potential produces the $(p, q)$-Laplace operator $\Delta_{p}+\Delta_{q}$, which generates a "double-phase energy" (according to the terminology of Marcellini and Mingione).

We briefly recall in what follows the roots of double-phase problems. To the best of our knowledge, problems of this type have been first considered by Ball $[6,7]$ in the context of problems with cavities in nonlinear elasticity.

Let $\Omega \subset \mathbb{R}^{N}(N \geqslant 2)$ be a bounded domain with smooth boundary. If $u: \Omega \rightarrow \mathbb{R}^{N}$ is the displacement and if $D u$ is the $N \times N$ matrix of the deformation gradient, then the total energy is defined by

$$
\begin{equation*}
E(u)=\int_{\Omega} f(x, D u(x)) d x, \tag{2}
\end{equation*}
$$

where $f=f(x, \xi): \Omega \times \mathbb{R}^{N \times N} \rightarrow \mathbb{R}$ is quasiconvex with respect to $\xi$. The simplest example considered by Ball is given by functions $f$ of the type

$$
f(\xi)=g(\xi)+h(\operatorname{det} \xi)
$$

where $\operatorname{det} \xi$ is the determinant of the $N \times N$ matrix $\xi$, and $g, h$ are nonnegative convex functions, which satisfy the growth conditions

$$
g(\xi) \geqslant c_{1}|\xi|^{p} ; \quad \lim _{t \rightarrow+\infty} h(t)=+\infty
$$

where $c_{1}$ is a positive constant and $1<p<N$. The condition $p<N$ is necessary to study the existence of equilibrium solutions with cavities, that is, minima of the integral (2) that are discontinuous at one point where a cavity forms. In fact, every $u$ with finite energy belongs to the Sobolev space $W^{1, p}\left(\Omega, \mathbb{R}^{N}\right)$, and thus it is a continuous function if $p>N$.

In accordance with these problems arising in nonlinear elasticity, Marcellini [17,18] considered continuous functions $f=f(x, u)$ with unbalanced growth that satisfy

$$
c_{1}|u|^{p} \leqslant|f(x, u)| \leqslant c_{2}\left(1+|u|^{q}\right) \quad \text { for all }(x, u) \in \Omega \times \mathbb{R},
$$

where $c_{1}, c_{2}$ are positive constants and $1 \leqslant p \leqslant q$. We also point out the contributions of Baroni, Colombo and Mingione [8,9] in the framework of non-autonomous functionals characterized by the fact that the energy density changes its ellipticity and growth properties according to the point.

These contributions are in relationship with the work of Zhikov [28], who described the behaviour of some phenomena arising in nonlinear elasticity. In fact, Zhikov intended to provide models for strongly anisotropic materials in the context of homogenization. For instance, Zhikov considered the "double-phase" energy functional defined by

$$
\begin{equation*}
\mathcal{P}_{p, q}(u):=\int_{\Omega}\left(|\nabla u|^{p}+a(x)|\nabla u|^{q}\right) d x, \quad 0 \leqslant a(x) \leqslant L, 1<p<q, \tag{3}
\end{equation*}
$$

where the modulating coefficient $a(x)$ dictates the geometry of the composite made by two differential materials, with hardening exponents $p$ and $q$, respectively. The functional $\mathcal{P}_{p, q}$ falls in the realm of the so-called functionals with nonstandard growth conditions of $(p, q)$-type, according to Marcellini's terminology. These are functionals of the type in (2), where the energy density satisfies

$$
|\xi|^{p} \leqslant f(x, \xi) \leqslant|\xi|^{q}+1, \quad 1 \leqslant p \leqslant q .
$$

General models with $(p, q)$-growth in the context of geometrically constrained problems have been recently studied by De Filippis [13]. This seems to be the first work dealing with $(p, q)$-conditions with manifold constraint. Refined regularity results are proved in [13], by using an approximation technique relying on estimates obtained through a careful use of difference quotients. Other recent works dealing with nonlinear problems with unbalanced growth (either isotropic or anisotropic) are the papers by Bahrouni, Rădulescu and Repovš [5], Cencelj, Rădulescu and Repovš [11], and Papageorgiou, Rădulescu and Repovš [19].

The differential operator defined in (1) and which is generated by a potential with variable growth was introduced by Azzollini et al. [2,3] in relationship with wide classes of nonlinear PDEs with a variational structure. We refer to Chorfi and Rădulescu [12] for the study of a related problem driven by this general differential operator. We also refer to the recent monograph $[20,26]$ for some of the abstract methods used in the present paper.

## 2. Functional setting and main results

We are concerned with the existence of nontrivial solutions of the following quasilinear Schrödinger problem with double-power nonlinearities:

$$
\begin{equation*}
-\operatorname{div}\left[\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u\right]+|u|^{\alpha-2} u=\lambda|u|^{s-2} u \quad \text { in } \Omega \subset \mathbb{R}^{N}(N \geqslant 2), \tag{4}
\end{equation*}
$$

where $\lambda$ is a positive parameter.
This equation was studied in [3] if $\Omega=\mathbb{R}^{N}$ and under the assumption that the reaction dominates the left-hand side of the problem. In fact, Azzollini, d'Avenia and Pomponio [3] proved that this equation has a nontrivial non-negative radially symmetric solution, provided that $1<p<q<\min \left\{N, p^{*}\right\}, 1<\alpha \leqslant p^{*} q^{\prime} / p^{\prime}$, and $\max \{q, \alpha\}<s<p^{*}$. A crucial tool in their arguments is a certain compactness property of a Sobolevtype space of radially symmetric functions into Lebesgue spaces. Our purpose in this paper is to consider the same equation but on exterior domains of the Euclidean space. We aim to prove related existence or non-existence results depending on the values of the parameter $\lambda$ and the competition between the left-hand side of Eq. (4) and its reaction. More precisely, we are first concerned with the following nonlinear eigenvalue problem

$$
\begin{cases}-\operatorname{div}\left[\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u\right]+|u|^{\alpha-2} u=\lambda|u|^{s-2} u & \text { in } \Omega  \tag{5}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

Here, we assume that $\lambda$ is a positive parameter and $\Omega \subset \mathbb{R}^{N}(N \geqslant 2)$ is the complement of a bounded domain with smooth boundary.

The existence of solutions of problem (5) was studied by Berestycki and Lions [10] in the case of the Laplace operator and without the presence of the nonlinear term $|u|^{\alpha-2} u$. In this case, the authors assumed
that the reaction is a double-power nonlinearity that has a subcritical behaviour at infinity and a supercritical growth near the origin.

In order to describe the main results of this paper, we start with the basic hypotheses.
Throughout this paper we assume that $\alpha, p, q$ are real numbers satisfying the following hypothesis:

$$
\begin{equation*}
1<p<q<N \quad \text { and } \quad 1<\alpha<\frac{p^{*} q^{\prime}}{p^{\prime}} \tag{6}
\end{equation*}
$$

We assume that the function $\phi: \mathbb{R}_{+} \rightarrow \mathbb{R}_{+}$is of class $C^{1}$ and has the following properties:
$\left(\phi_{1}\right) \phi(0)=0$;
$\left(\phi_{2}\right)$ there exists $c>0$ such that $\phi(t) \geqslant c t^{p / 2}$ if $t \geqslant 1$ and $\phi(t) \geqslant c t^{q / 2}$ if $0 \leqslant t \leqslant 1$;
$\left(\phi_{3}\right)$ there exists $C>0$ such that $\phi(t) \leqslant C t^{p / 2}$ if $t \geqslant 1$ and $\phi(t) \leqslant C t^{q / 2}$ if $0 \leqslant t \leqslant 1$;
$\left(\phi_{4}\right)$ there exists $0<\mu<1$ such that $2 t \phi^{\prime}(t) \leqslant s \mu \phi(t)$ for all $t \geqslant 0$;
$\left(\phi_{5}\right)$ the mapping $t \mapsto \phi\left(t^{2}\right)$ is strictly convex.
Since our hypotheses allow that $\phi^{\prime}$ approaches 0 , problem (5) is degenerate and no ellipticity condition is assumed.

For all $1 \leqslant r \leqslant \infty$, we denote by $\|\cdot\|_{r}$ the norm on the Lebesgue space $L^{r}(\Omega)$.

Definition 1. Let $\Omega \subseteq \mathbb{R}^{N}$ be an open set. We define the function space $L^{p}(\Omega)+L^{q}(\Omega)$ as the completion of $C_{c}^{\infty}(\Omega)$ in the norm

$$
\|u\|_{L^{p}+L^{q}}:=\inf \left\{\|v\|_{p}+\|w\|_{q} ; v \in L^{p}(\Omega), w \in L^{q}(\Omega), u=v+w\right\} .
$$

We set

$$
\|u\|_{p, q}:=\|u\|_{L^{p}(\Omega)+L^{q}(\Omega)} .
$$

The space $L^{p}(\Omega)+L^{q}(\Omega)$ is an Orlicz space and has been intensively studied by Badiale, Pisani and Rolando [4, Sect. 2]. This space is a reflexive Banach space, see [4, Corollary 2.11]. We point out that the space $L^{p}(\Omega)+L^{q}(\Omega)$ is of interest only either $p<q$ or $|\Omega|=+\infty$. Indeed, if $p=q$ or $|\Omega|<+\infty$, then $L^{q}(\Omega) \subseteq L^{p}(\Omega)$, hence $L^{p}(\Omega)+L^{q}(\Omega)=L^{p}(\Omega)$.

A key role in our arguments is played by the Banach space

$$
\mathcal{B}:={\overline{C_{c}^{\infty}(\Omega)}}^{\|\cdot\|},
$$

where

$$
\|u\|:=\|\nabla u\|_{p, q}+\|u\|_{\alpha} .
$$

As established in Propositions 2.4 and 2.5 of [3], $\mathcal{B}$ is a reflexive Banach space. Moreover, if $p^{\prime}<p^{*} q^{\prime}$ then for every $1<\alpha \leqslant p^{*} q^{\prime} / p^{\prime}$, the space $\mathcal{B}$ is continuously embedded into $L^{p^{*}}(\Omega)$; see [3, Theorem 2.6] for more details. We point out that the loss of compactness of the Orlicz embeddings in the case of unbounded domains implies refined variational techniques. Some of the papers dealing with problems with lack of compactness on unbounded domains use particular function spaces where the compactness is preserved, such as spaces of radially symmetric functions. Such a situation occurs in [3], where the main existence property is obtained via a compact embedding. We recall that even if the domain is unbounded, standard compact embeddings remain true, for instance if $\Omega$ is "thin at infinity", in the sense that

$$
\lim _{R \rightarrow \infty} \sup \left\{|\Omega \cap B(x, 1)| ; x \in \mathbb{R}^{N},|x|=R\right\}=0
$$

Such a situation does not hold in our case. Indeed, since $\Omega$ is an exterior domain, then it looks like the whole space $\mathbb{R}^{N}$ at infinity and, in particular, it is not a thin domain.

Definition 2. A solution of problem (5) is a function $u \in \mathcal{B} \backslash\{0\}$ such that for all $v \in \mathcal{B}$

$$
\int_{\Omega}\left(\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u \nabla v+|u|^{\alpha-2} u v-\lambda|u|^{s-2} u v\right) d x=0 .
$$

The real number $\lambda$ for which problem (5) has a nontrivial solution is an eigenvalue and the corresponding $u \in \mathcal{B} \backslash\{0\}$ is an eigenfunction of the problem. These terms are in accordance with the related notions introduced by Fučik, Nečas, Souček and Souček [16, p. 117] in the abstract framework of nonlinear operators. Indeed, if we set

$$
S u:=\frac{1}{2} \int_{\Omega} \phi\left(|\nabla u|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega}|u|^{\alpha} d x \quad \text { and } \quad T(u):=\frac{1}{s} \int_{\Omega}|u|^{s} d x
$$

then $\lambda$ is an eigenvalue for the pair $(S, T)$ if and only if there exists a corresponding eigenfunction, namely a solution of problem (5) as described by Definition 2.

We first prove that problem (5) has a solution for any $\lambda>0$, provided that the reaction "dominates" the growth in the left-hand side. More precisely, we have the following existence result.

Theorem 3. Assume that hypotheses (6), ( $\phi_{1}$ )-( $\phi_{5}$ ) are fulfilled, and $\max \{q, \alpha\}<s<p^{*}$. Then the following properties are true:
(a) problem (5) has a nonnegative solution $U$ for all $\lambda>0$;
(b) $U \in C^{1, \mu}\left(\Omega \cap B_{R}(0)\right)$ with $\mu=\mu(R) \in(0,1)$;
(c) $U>0$ in $\Omega$.

Next, we are concerned with the following nonlinear problem with variable potential and lack of compactness

$$
\begin{cases}-\operatorname{div}\left[\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u\right]+|u|^{\alpha-2} u=\lambda a(x)|u|^{s-2} u & \text { in } \Omega  \tag{7}\\ u=0 & \text { on } \partial \Omega .\end{cases}
$$

Accordingly, a solution of problem (7) is a function $u \in \mathcal{B} \backslash\{0\}$ such that for all $v \in \mathcal{B}$

$$
\int_{\Omega}\left(\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u \nabla v+|u|^{\alpha-2} u v-\lambda a(x)|u|^{s-2} u v\right) d x=0 .
$$

Hypothesis (6) is now replaced by

$$
\begin{equation*}
\max \{q, s\}<\alpha<p^{*} \tag{8}
\end{equation*}
$$

We assume that the potential $a \geqslant 0$ is positive on a subset of $\Omega$ of positive measure and

$$
\begin{equation*}
a^{\alpha /(\alpha-s)} \in L^{1}(\Omega) \tag{9}
\end{equation*}
$$

The second main result of this paper establishes an existence and non-existence property if the reaction of problem (7) is dominated by the left-hand side. In this case, solutions exist only for high perturbations of the right-hand side.

Theorem 4. Assume that hypotheses (6), (8), (9), and ( $\phi_{1}$ )-( $\phi_{5}$ ) are fulfilled. Then there exists $\Lambda>0$ such that the following properties are true:
(a) problem (7) does not have any solution for all $0<\lambda<\Lambda$;
(b) problem (7) has a positive solution $U$ for all $\lambda \geqslant \Lambda$. Moreover, $U \in C^{1, \mu}\left(\Omega \cap B_{R}(0)\right)$ with $\mu=\mu(R) \in$ $(0,1)$.

These results remain true if we replace the power-type nonlinearities with general nonlinearities. For instance, the reaction $|u|^{s-2} u$ in the statement of Theorem 3 corresponding to problem (5), can be replaced by a Carathéodory function $f: \mathbb{R}^{N} \times \mathbb{R} \rightarrow \mathbb{R}$ with the following properties:
$\left(f_{1}\right) f(x, u)=o\left(u^{\alpha-1}\right)$ as $u \rightarrow 0^{+}$, uniformly for a.e. $x \in \mathbb{R}^{N}$;
$\left(f_{2}\right) f(x, u)=O\left(u^{s-1}\right)$ as $u \rightarrow+\infty$, uniformly for a.e. $x \in \mathbb{R}^{N}$.
The above results extend some related properties established by Filippucci, Pucci and Rădulescu [15] in the framework of the $p$-Laplace operator. We also refer to Chorfi and Rădulescu [12] who studied a related problem driven by the same differential operator and if $\Omega=\mathbb{R}^{N}$.

## 3. Proof of Theorem 3

We point out that a related property is proved by Azzollini, d'Avenia and Pomponio [3] if $\Omega=\mathbb{R}^{N}$. However, Theorem 1.3 in [3] establishes the existence of a radially symmetric solution and the proof strongly relies on the compact embedding of a Sobolev-type space of functions with radially symmetry into certain Lebesgue spaces. In our case, since $\Omega$ is unbounded but without any symmetry properties, we are not looking for radially symmetric solutions.

The energy functional associated to problem (5) is $\mathcal{E}: \mathcal{B} \rightarrow \mathbb{R}$ defined by

$$
\mathcal{E}(u):=\frac{1}{2} \int_{\Omega} \phi\left(|\nabla u|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega}|u|^{\alpha} d x-\frac{\lambda}{s} \int_{\Omega}|u|^{s} d x .
$$

By [3, Proposition 3.1], $\mathcal{E}$ is well-defined and of class $C^{1}$. Moreover, for all $u, v \in \mathcal{B}$ its Gâteaux directional derivative is given by

$$
\left\langle\mathcal{E}^{\prime}(u), v\right\rangle=\int_{\Omega}\left(\phi^{\prime}\left(|\nabla u|^{2}\right) \nabla u \nabla v+|u|^{\alpha-2} u v-\lambda|u|^{s-2} u v\right) d x
$$

We first claim that

$$
\begin{equation*}
\text { there exists small } r>0 \text { such that } \inf _{\|u\|=r} \mathcal{E}(u)>0 \text {. } \tag{10}
\end{equation*}
$$

By $\left(\phi_{1}\right)$ we have for all $u \in \mathcal{B}$

$$
\begin{equation*}
\mathcal{E}(u) \geqslant \frac{c}{2} \int_{|\nabla u| \leqslant 1}|\nabla u|^{q} d x+\frac{c}{2} \int_{|\nabla u|>1}|\nabla u|^{p} d x+\frac{1}{\alpha} \int_{\Omega}|u|^{\alpha} d x-\frac{\lambda}{s} \int_{\Omega}|u|^{s} d x . \tag{11}
\end{equation*}
$$

By [3, Theorem 2.6] and our hypothesis $\max \{q, \alpha\}<s<p^{*}$, it follows that $\mathcal{B}$ is continuously embedded into $L^{s}(\Omega)$. So, there exists $c_{1}>0$ such that

$$
\|u\|_{s} \leqslant c_{1}\|u\| \quad \text { for all } u \in \mathcal{B}
$$

Returning to (11) we obtain for all $u \in \mathcal{B}$

$$
\begin{equation*}
\mathcal{E}(u) \geqslant \frac{c}{2} \int_{|\nabla u| \leqslant 1}|\nabla u|^{q} d x+\frac{c}{2} \int_{|\nabla u|>1}|\nabla u|^{p} d x+\frac{1}{\alpha}\|u\|_{\alpha}^{\alpha}-\frac{\lambda c_{1}}{s}\|u\|^{s} . \tag{12}
\end{equation*}
$$

Fix $r \in(0,1)$. By (6) and (12) we deduce that there are positive constants $c_{2}, c_{3}$ and $c_{4}$ such that for all $u \in \mathcal{B}$ with $\|u\|=1$

$$
\begin{align*}
\mathcal{E}(u) & \geqslant c_{2}\left(\|\nabla u\|_{p, q}^{q}+\|u\|_{\alpha}^{\alpha}\right)-c_{3}\|u\|^{s} \\
& \geqslant c_{4}\left(\|u\|^{\max \{q, \alpha\}}-\|u\|^{s}\right) . \tag{13}
\end{align*}
$$

Since $\max \{q, \alpha\}<s<p^{*}$, relation (13) shows that there exists $c_{5}>0$ such that

$$
\begin{equation*}
\mathcal{E}(u) \geqslant c_{5} \quad \text { for all } u \in \mathcal{B} \text { with }\|u\|=r, \tag{14}
\end{equation*}
$$

which proves (10).
Next, we claim that

$$
\begin{equation*}
\liminf _{R \rightarrow+\infty} \mathcal{E}(u)=-\infty \tag{15}
\end{equation*}
$$

Indeed, fix $\psi \in \mathcal{B} \backslash\{0\}$ and $t>0$. Thus, by ( $\phi_{3}$ ),

$$
\begin{aligned}
\mathcal{E}(t \psi) & \leqslant \frac{C}{2}\left(t^{q} \int_{t|\nabla \psi| \leqslant 1}|\nabla \psi|^{q} d x+t^{p} \int_{t|\nabla \psi|>1}|\nabla \psi|^{p} d x\right)+\frac{t^{\alpha}}{\alpha} \int_{\Omega}|\psi|^{\alpha} d x-\frac{\lambda t^{s}}{s} \int_{\Omega}|\psi|^{s} d x \\
& =A_{1} t^{q}+A_{2} t^{p}+A_{3} t^{\alpha}-A_{4} t^{s}\left(\text { with } A_{1}, A_{2} \geqslant 0, A_{3}, A_{4}>0\right) \rightarrow-\infty \text { as } t \rightarrow+\infty,
\end{aligned}
$$

by our hypothesis. This proves (15).
By relations (10), (15) and using the mountain pass theorem, we find $\left(u_{n}\right) \subset \mathcal{B}$ such that

$$
\begin{equation*}
\mathcal{E}\left(u_{n}\right) \rightarrow c_{0} \quad \text { and } \quad \mathcal{E}^{\prime}\left(u_{n}\right) \rightarrow 0 \text { in } \mathcal{B}^{*} \quad \text { as } n \rightarrow \infty . \tag{16}
\end{equation*}
$$

Here, $c_{0}:=\inf _{\gamma \in \mathcal{C}} \max _{t \in[0,1]} \mathcal{E}(\gamma(t))>0$, where

$$
\mathcal{C}:=\left\{\gamma:[0,1] \rightarrow \mathcal{B} ; \gamma \text { is continuous, } \gamma(0)=0, \gamma(1)=t_{0} \psi\right\},
$$

for some fixed $t_{0}>0$ such that $t_{0}\|\psi\|>c_{5}$, where $c_{5}$ is defined in (14).
Combining ( $\phi_{4}$ ) and ( $\phi_{5}$ ) we deduce that

$$
\begin{equation*}
\phi(t) \leqslant 2 t \phi^{\prime}(t) \leqslant s \mu \phi(t) \quad \text { for all } t>0, \tag{17}
\end{equation*}
$$

hence $s \mu>1$. Thus, by (17), $\phi$ is increasing. It follows that

$$
\mathcal{E}(|v|) \leqslant \mathcal{E}(v) \quad \text { for all } v \in \mathcal{B} .
$$

We deduce that we can assume that $u_{n} \geqslant 0$ in (16).
If ( $u_{n}$ ) satisfies (16) then

$$
\begin{equation*}
\mathcal{E}\left(u_{n}\right)-\frac{1}{s}\left\langle\mathcal{E}^{\prime}\left(u_{n}\right), u_{n}\right\rangle=O(1)+o\left(\left\|u_{n}\right\|\right) \quad \text { as } n \rightarrow \infty \tag{18}
\end{equation*}
$$

But, by ( $\phi_{4}$ )

$$
\begin{align*}
\mathcal{E}\left(u_{n}\right)-\frac{1}{s}\left\langle\mathcal{E}^{\prime}\left(u_{n}\right), u_{n}\right\rangle & =\int_{\Omega}\left(\frac{1}{2} \phi\left(\left|\nabla u_{n}\right|^{2}\right)-\frac{1}{s} \phi^{\prime}\left(\left|\nabla u_{n}\right|^{2}\right)\left|\nabla u_{n}\right|^{2}\right) d x \\
& +\left(\frac{1}{\alpha}-\frac{1}{s}\right) \int_{\Omega}\left|u_{n}\right|^{\alpha} d x  \tag{19}\\
& \geqslant \frac{1-\mu}{2} \int_{\Omega} \phi\left(\left|\nabla u_{n}\right|^{2}\right) d x+\left(\frac{1}{\alpha}-\frac{1}{s}\right) \int_{\Omega}\left|u_{n}\right|^{\alpha} d x \\
& =c_{6} \int_{\Omega} \phi\left(\left|\nabla u_{n}\right|^{2}\right) d x+c_{7}\left\|u_{n}\right\|_{\alpha}^{\alpha}
\end{align*}
$$

where $c_{6}, c_{7}>0$.
Next, with an argument similar as the same developed in the first part of this proof, relation (19) implies that for some $c_{8}>0$

$$
\begin{equation*}
\mathcal{E}\left(u_{n}\right)-\frac{1}{s}\left\langle\mathcal{E}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \geqslant c_{8}\left(\left\|\nabla u_{n}\right\|_{p, q}^{q}+\left\|u_{n}\right\|^{\alpha}\right) \quad \text { for all } n \geqslant 1 . \tag{20}
\end{equation*}
$$

Combining (18) and (20) we deduce that

$$
\left\|\nabla u_{n}\right\|_{p, q}^{q}+\left\|u_{n}\right\|^{\alpha} \leqslant O(1)+o\left(\left\|u_{n}\right\|\right) \quad \text { as } n \rightarrow \infty,
$$

which shows that $\left(u_{n}\right)$ is bounded in $\mathcal{B}$.
Until now we have proved that the Palais-Smale sequence $\left(u_{n}\right)$ of $\mathcal{E}$ is bounded. Thus, there exists $U \in \mathcal{B}$ such that, up to a subsequence,

$$
u_{n} \rightharpoonup U \geqslant 0 \quad \text { in } \mathcal{B}
$$

and

$$
u_{n} \rightarrow U \quad \text { in } L^{s}(\Omega) \text { and } L^{\alpha}(\Omega) .
$$

We prove in what follows that $U$ is a solution of problem (5). For this purpose we fix $v \in C_{c}^{\infty}(\Omega)$ and we set $\omega:=\operatorname{supp} v$. Define the functional

$$
\mathcal{E}_{0}(u):=\frac{1}{2} \int_{\omega} \phi\left(|\nabla u|^{2}\right) d x+\frac{1}{\alpha} \int_{\omega}|u|^{\alpha} d x .
$$

By $\left(\phi_{5}\right)$ it follows that $\mathcal{E}_{0}$ is convex. Since it is also continuous, it follows that $\mathcal{E}_{0}$ is weakly lower semicontinuous. By convexity we have

$$
\mathcal{E}_{0}\left(u_{n}\right) \leqslant \mathcal{E}_{0}(U)+\left\langle\mathcal{E}_{0}^{\prime}\left(u_{n}\right), u_{n}-U\right\rangle .
$$

By (16) we deduce that

$$
\limsup _{n \rightarrow \infty} \mathcal{E}_{0}\left(u_{n}\right) \leqslant \mathcal{E}_{0}(U)
$$

Using now the weakly lower semicontinuity of $\mathcal{E}_{0}$ we conclude that

$$
\limsup _{n \rightarrow \infty} \mathcal{E}_{0}\left(u_{n}\right)=\mathcal{E}_{0}(U) .
$$

Next, with the same arguments as in [3, p. 210], it follows that

$$
\nabla u_{n} \rightarrow \nabla U \quad \text { in } L^{p}(\Omega)+L^{q}(\Omega) .
$$

Using (16) and passing to the limit as $n \rightarrow \infty$ we deduce that

$$
\int_{\omega} \phi^{\prime}\left(|\nabla U|^{2}\right) \nabla U \nabla v d x+\int_{\omega}|U|^{\alpha-2} U v d x-\lambda \int_{\omega}|U|^{s-2} U v d x=0 .
$$

By density, this identity is valid for any $v \in \mathcal{B}$. Thus, $U$ is a solution of problem (5).
We prove in what follows that $U \neq 0$. Indeed, if not, it follows that

$$
u_{n} \rightarrow 0 \quad \text { in } L^{s}(\Omega) \text { and } L^{\alpha}(\Omega) .
$$

Thus, by (16)

$$
\begin{aligned}
\frac{c_{0}}{2} & \leqslant \mathcal{E}\left(u_{n}\right)-\frac{1}{2}\left\langle\mathcal{E}^{\prime}\left(u_{n}\right), u_{n}\right\rangle \\
& =\frac{1}{2} \int_{\Omega}\left(\phi\left(\left|\nabla u_{n}\right|^{2}\right)-\phi^{\prime}\left(\left|\nabla u_{n}\right|^{2}\right)\left|\nabla u_{n}\right|^{2}\right) d x+\left(\frac{1}{\alpha}-\frac{1}{2}\right) \int_{\Omega}\left|u_{n}\right|^{\alpha} d x \\
& +\lambda\left(\frac{1}{2}-\frac{1}{s}\right) \int_{\Omega}\left|u_{n}\right|^{s} d x
\end{aligned}
$$

By $\left(\phi_{5}\right)$ it follows that $\phi\left(t^{2}\right)-\phi^{\prime}\left(t^{2}\right) t^{2} \leqslant 0$, hence

$$
0<\frac{c_{0}}{2} \leqslant\left(\frac{1}{\alpha}-\frac{1}{2}\right) \int_{\Omega}\left|u_{n}\right|^{\alpha} d x+\lambda\left(\frac{1}{2}-\frac{1}{s}\right) \int_{\Omega}\left|u_{n}\right|^{s} d x .
$$

Passing to the limit as $n \rightarrow \infty$ we get a contradiction.
We conclude that $U \neq 0$ and $U \geqslant 0$.
(b) By Theorem 1(ii) of Pucci and Servadei [24], which is based on the Moser iteration, we first deduce that $U \in L_{l o c}^{\infty}(\Omega)$. Next, using the corollary of DiBenedetto [14, p. 830], we conclude that $U \in C^{1, \mu}\left(\Omega \cap B_{R}(0)\right)$ with $\mu=\mu(R) \in(0,1)$. A related argument was applied in the proof of Theorem 1 in Yu [27].
(c) The function $U \geqslant 0$ satisfies

$$
-\operatorname{div}\left[\phi^{\prime}\left(|\nabla U|^{2}\right) \nabla U\right]+|U|^{\alpha-2} U \geqslant 0 \text { in } \Omega .
$$

We recall that the generalized maximum principle of Pucci and Serrin [21-23] applied to general canonical divergence structure inequalities of the type

$$
-\operatorname{div}(A(|\nabla u|) \nabla u)+f(u) \geqslant 0 \quad \text { in } \Omega,
$$

where the function $A=A(t)$ and the nonlinearity $f$ satisfy the following conditions:
(A1) $A$ is continuous in $\mathbb{R}^{+}$;
(A2) the mapping $t \mapsto t A(t)$ is strictly increasing in $\mathbb{R}^{+}$and $t A(t) \rightarrow 0$ as $t \rightarrow 0^{+}$;
(F1) $f \in C\left(\mathbb{R}_{0}^{+}\right)$;
(F2) $f(0)=0$ and $f$ is non-decreasing on some interval $(0, \delta), \delta>0$.
In our case, $f(u)=|u|^{\alpha-2} u$ satisfies (F1) and (F2). We have $A(t)=\phi^{\prime}\left(t^{2}\right)$ and $t \phi^{\prime}\left(t^{2}\right)$ is strictly increasing by our hypothesis $\left(\phi_{5}\right)$. We also observe that $\left(\phi_{4}\right)$ and $\left(\phi_{3}\right)$ imply for all $t \in(0,1)$

$$
0<t \phi^{\prime}\left(t^{2}\right) \leqslant \frac{s \mu}{2} \frac{\phi\left(t^{2}\right)}{t} \leqslant \frac{s \mu}{2} t^{q-1} \rightarrow 0 \text { as } t \rightarrow 0^{+}
$$

So, by the Pucci-Serrin maximum principle, we conclude that the non-negative solution $U$ is positive in $\Omega$.

## 4. Proof of Theorem 4

The energy functional associated to problem (7) is $\mathcal{F}: \mathcal{B} \rightarrow \mathbb{R}$ defined by

$$
\mathcal{F}(u):=\frac{1}{2} \int_{\Omega} \phi\left(|\nabla u|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega}|u|^{\alpha} d x-\frac{\lambda}{s} \int_{\Omega} a(x)|u|^{s} d x
$$

We first establish that $\mathcal{F}$ is coercive. Indeed, by ( $\phi_{2}$ ) we have for all $u \in \mathcal{B}$

$$
\begin{equation*}
\mathcal{F}(u) \geqslant \frac{c}{2} \int_{|\nabla u|>1}|\nabla u|^{p} d x+\frac{c}{2} \int_{|\nabla u| \leqslant 1}|\nabla u|^{q} d x+\frac{1}{\alpha}\|u\|_{\alpha}^{\alpha}-\frac{\lambda}{s} \int_{\Omega} a(x)|u|^{s} d x . \tag{21}
\end{equation*}
$$

By Hölder's inequality and hypothesis (9) we obtain

$$
\begin{equation*}
\int_{\Omega} a(x)|u|^{s} d x \leqslant\|a\|_{\alpha /(\alpha-s)} \cdot\left(\int_{\Omega}|u|^{\alpha} d x\right)^{s / \alpha}=C_{1}\|u\|_{\alpha}^{s} \tag{22}
\end{equation*}
$$

where $C_{1}=C_{1}(a, \alpha, s, \Omega)$.
Since $\alpha>s$, relations (21) and (22) yield

$$
\mathcal{F}(u) \geqslant \frac{c}{2}\|\nabla u\|_{p, q}^{p}+\frac{1}{\alpha}\|u\|_{\alpha}^{\alpha}-C_{1}\|u\|_{\alpha}^{s} \rightarrow+\infty \quad \text { as }\|u\| \rightarrow \infty,
$$

hence $\mathcal{F}$ is coercive and bounded from below.
We prove in what follows that problem (7) does not have any solution, provided that $\lambda>0$ is sufficiently small. Indeed, we observe that if $u$ solves (7) then

$$
\int_{\Omega} \phi^{\prime}\left(|\nabla u|^{2}\right)|\nabla u|^{2} d x+\int_{\Omega}|u|^{\alpha} d x=\lambda \int_{\Omega} a(x)|u|^{s} d x
$$

We now estimate the right-hand side of this equality. By hypotheses (8) (9) we have

$$
\begin{aligned}
\lambda \int_{\Omega} a(x)|u|^{s} d x & \leqslant \lambda^{\alpha /(\alpha-s)} \frac{\alpha-s}{\alpha} \int_{\Omega} a(x)^{\alpha /(\alpha-s)} d x+\frac{s}{\alpha} \int_{\Omega}|u|^{\alpha} d x \\
& =C(s, a, \alpha) \lambda^{\alpha /(\alpha-s)}+\frac{s}{\alpha} \int_{\Omega}|u|^{\alpha} d x
\end{aligned}
$$

We deduce that if $u$ is a solution of problem (7) then

$$
\begin{aligned}
0 & \leqslant \int_{\Omega} \phi^{\prime}\left(|\nabla u|^{2}\right)|\nabla u|^{2} d x \leqslant C(s, a, \alpha) \lambda^{\alpha /(\alpha-s)}+\left(\frac{s}{\alpha}-1\right) \int_{\Omega}|u|^{\alpha} d x \\
& <C(s, a, \alpha) \lambda^{\alpha /(\alpha-s)}
\end{aligned}
$$

by (8).
In conclusion, problem (7) does not have any solution, provided that $\lambda>0$ is small enough. Let

$$
\lambda_{*}:=\sup \{\lambda>0 ; \text { problem (7) does not have a solution }\}>0 .
$$

The above arguments show that (7) does not have a solution for all $\lambda<\lambda_{*}$.
In order to obtain sufficient conditions for the existence of solutions, we consider the minimization problem

$$
m:=\inf _{u \in \mathcal{B}} \mathcal{F}(u) \in \mathbb{R} .
$$

Let $\left(u_{n}\right) \subset \mathcal{B}$ be a minimizing sequence of $\mathcal{F}$. Since $\mathcal{F}\left(\left|u_{n}\right|\right) \leqslant \mathcal{F}\left(u_{n}\right)$, we can assume that $u_{n} \geqslant 0$. Moreover, $\left(u_{n}\right)$ is bounded so, up to a subsequence, we can assume that

$$
u_{n} \rightharpoonup U \geqslant 0 \quad \text { in } \mathcal{B} .
$$

Hypothesis (8) implies that $\mathcal{B}$ is compactly embedded into the weighted Lebesgue space $L^{s}(\Omega ; a)$. So, by weak lower semicontinuity and compactness of the embedding, we have

$$
\frac{1}{2} \int_{\Omega} \phi\left(|\nabla U|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega} U^{\alpha} d x \leqslant \liminf _{n \rightarrow \infty}\left(\frac{1}{2} \int_{\Omega} \phi\left(\left|\nabla u_{n}\right|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega} u_{n}^{\alpha} d x\right)
$$

and

$$
\int_{\Omega} a(x) u_{n}^{s} d x \rightarrow \int_{\Omega} a(x) U^{s} d x \quad \text { as } n \rightarrow \infty
$$

It follows that $U \geqslant 0$ is a minimizer of $\mathcal{F}$, that is, $\mathcal{F}(U)=m$.
We now prove that $U$ is a solution of problem (7), provided that $\lambda$ is big enough. For this purpose, consider the minimization problem

$$
\begin{equation*}
m_{0}:=\inf _{w \in \mathcal{B}}\left\{\frac{1}{2} \int_{\Omega} \phi\left(|\nabla w|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega}|w|^{\alpha} d x ; \frac{1}{s} \int_{\Omega} a(x)|w|^{s} d x=1\right\} \tag{23}
\end{equation*}
$$

If $\left(w_{n}\right) \subset \mathcal{B}$ is a minimizing sequence, then $\left(w_{n}\right)$ is bounded. So, up to a subsequence, we can assume that

$$
\begin{gathered}
w_{n} \rightharpoonup w \quad \text { in } \mathcal{B} \\
w_{n} \rightarrow w \quad \text { in } L^{s}(\Omega ; a) .
\end{gathered}
$$

It follows that $w$ is a solution of (23), hence $\mathcal{F}(w)=m_{0}-\lambda$. We deduce that problem (7) has a solution for all $\lambda>m_{0}$.

We set

$$
\lambda^{*}:=\inf \{\lambda>0 ; \text { problem (7) has a solution }\} .
$$

Then $\lambda^{*} \geqslant \lambda_{*}$.

Next, we prove that (7) has a solution for all $\lambda>\lambda^{*}$. Indeed, if we fix $\lambda>\lambda^{*}$, then the definition of $\lambda^{*}$ yields some $\lambda^{*}<\underline{\lambda}<\lambda$ such that problem (7) has a solution $\underline{U}$ corresponding to $\underline{\lambda}$. Then $\underline{U}$ is a subsolution of (7). It remains to prove that problem (7) has a supersolution $\bar{U}$ such that $\bar{U} \geqslant \underline{U}$. For this purpose we consider the new minimization problem

$$
\begin{equation*}
\inf _{v \in \mathcal{B}}\left\{\frac{1}{2} \int_{\Omega} \phi\left(|\nabla v|^{2}\right) d x+\frac{1}{\alpha} \int_{\Omega}|v|^{\alpha} d x-\frac{\lambda}{s} \int_{\Omega} a(x)|v|^{s} d x ; v \geqslant \underline{U}\right\} \tag{24}
\end{equation*}
$$

Using the same arguments as above we deduce that the constrained minimization problem (24) has a solution $\bar{U} \geqslant \underline{U}$. We conclude that (7) has a solution for all $\lambda>\lambda^{*}$.

The definition of $\lambda^{*}$ shows that problem (7) has no solution if $0<\lambda<\lambda^{*}$. Since $\lambda^{*} \geqslant \lambda_{*}$, we conclude that

$$
\lambda^{*}=\lambda_{*}=: \Lambda
$$

Until now we know that (7) has no solution if $0<\lambda<\Lambda$ but it has at least one non-negative solution $U$ for all $\lambda \geqslant \Lambda$. We now prove that problem (7) has a non-negative solution if $\lambda=\Lambda$. Indeed, let $\left(\lambda_{n}\right)$ be a sequence of real numbers such that $\lambda_{n} \downarrow \Lambda$ as $n \rightarrow \infty$. Let $U_{n} \geqslant 0$ be a solution of (7) corresponding to $\lambda_{n}$. Since $\left(U_{n}\right) \subset \mathcal{B}$ is bounded, we can assume, passing eventually to a subsequence, that

$$
\begin{gather*}
U_{n} \rightharpoonup U_{\Lambda} \quad \text { in } \mathcal{B}  \tag{25}\\
U_{n} \rightarrow U_{\Lambda} \quad \text { in } L^{s}(\Omega ; a)  \tag{26}\\
U_{n} \rightarrow U_{\Lambda} \quad \text { a.e. } \Omega \tag{27}
\end{gather*}
$$

Since $U_{n}$ solves (7) for $\lambda=\lambda_{n}$, it follows that for all $v \in \mathcal{B}$

$$
\begin{equation*}
\int_{\Omega} \phi^{\prime}\left(\left|\nabla U_{n}\right|^{2}\right) \nabla U_{n} \nabla v d x+\int_{\Omega} U_{n}^{\alpha-2} U_{n} v d x=\lambda_{n} \int_{\Omega} U_{n}^{s-2} U_{n} v d x=0 \text { for all } n \geqslant 1 \tag{28}
\end{equation*}
$$

Taking $n \rightarrow \infty$ in (28) and using (25)-(27), we deduce that $U_{\Lambda} \geqslant 0$ is a solution of problem (7) for $\lambda=\Lambda$. We conclude that problem (7) has a solution $U \geqslant 0$ for every $\lambda \geqslant \Lambda$.

Next, as in the proof of Theorem 3(b) and using Theorem 1(ii) of Pucci and Servadei [24] in combination with the Moser iteration, we deduce that $U \in L_{l o c}^{\infty}(\Omega)$. This regularity property implies that $U \in C^{1, \mu}(\Omega \cap$ $\left.B_{R}(0)\right)$, where $\mu=\mu(R) \in(0,1)$ Applying the generalized Pucci-Serrin maximum principle, as in the proof of Theorem 3(c), we conclude that $U>0$ in $\Omega$.

## Final comments

We consider that an interesting research direction with multiple applications concerns the study of nonlinear problems described by the nonlocal term

$$
M\left(\int \phi\left(|\nabla u|^{2}\right)|\nabla u|^{2}\right)
$$

where $\phi$ satisfies hypotheses $\left(\phi_{1}\right)-\left(\phi_{5}\right)$. Pioneering results have been established by Pucci at all. [1,25] in the framework of Kirchhoff problems involving nonlocal operators associated to the standard differential operators.

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